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## THE QUAD SCHEME: A New Method for Phase Space Integration

by

F. M. Mueller, J. W. Garland,  
M. H. Cohen, and K. H. Bennemann

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Solid State Science Division

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ABSTRACT

The QUAD scheme is a general computational method for using the results of band calculations to compute integrals over momentum space. A local quadratic expansion is used to fit the electronic band structure over a sufficiently small cubic region of the Brillouin zone. By means of this expansion, the contribution of the region to such quantities as the density of states is found by Monte Carlo sampling. The total density of states is formed by adding the contributions from all the small cubic regions. The QUAD procedure should be applicable to many other calculations. Fortran IV programs are included for the problem of density of states and the combined interpolation scheme. Parameters appropriate to relativistic fcc platinum bands are presented.

I. INTRODUCTION

Large-scale electronic computers have made possible detailed comparisons between theoretical calculations of electronic properties of crystalline systems and experimental results. The problem in making such comparisons has been in actually performing summations over momentum space. The summations involve evaluation of singular three-dimensional integrals. In the past the problem could be treated in two ways. Systems could be chosen which were amenable to simple approximations to both the eigenvalues and eigenvectors; the necessary integrals could then be performed analytically. Alternatively, experiments could be chosen which depended only on limited regions of the Brillouin zone (BZ), perhaps regions near the Fermi surface or near specific critical points such as symmetry points.

Such approximations and restrictions are quite appropriate for simple metals and semiconductors. However, the transition metals, rare earths, and actinides cannot be so approximated. Every eigenvector and matrix element in these systems has a strong  $\mathbf{k}$  dependence because of hybridization between either the s and d electrons in the transition metals

and rare earths, or the s, d, and f electrons in the actinides. Mueller and Phillips<sup>1</sup> have shown that these materials have oscillator strengths which vary from 7.0 eV to 0.0 eV over a range in the BZ as short as 10% of a zone diameter. The energy bands of these materials are both numerous and convoluted. Even calculation of a simple property such as the density of states, which may be viewed as a "process" in which all the matrix elements are strictly constant, is difficult.

Recently a general technique, the QUAD scheme, has been developed for calculating the electronic properties of regular solids. The QUAD scheme divides an appropriate sector of the BZ into a sequence of small cubes. Eigenvalues and eigenvectors are found at a series of points within one of these cubes. These points are used to form expansion coefficients to full quadratic order by a least-squares procedure. The density of states over the cell is formed by using the expansion functions and by Monte Carlo sampling.

Such a sampling procedure for finding the density of states is approximate. One could, in principle, perform an exact calculation by finding the fractional volume of the cell enclosed by constant energy contours. Gilat and Raubenheimer<sup>2</sup> used this technique with a linear expansion to treat phonon spectra. A linear expansion is adequate for phonons when the bands are few in number and have little structure. Volume integration may be carried out quite easily in this linear approximation.

Electron bands of transition metals, however, are more numerous and much more convoluted than are phonon bands. The range of validity of the linear-expansion approximation for electron bands of transition metals is less than in phonon bands. Tests indicate that, if the expansion is to be accurate to 0.1% of the band width, a linear expansion breaks down about five times faster in the d-band region of the electronic band structure of copper as in the phonon bands of copper. For the same percent accuracy of approximation a net five times finer would be needed in the electron-band problem as in the phonon problem. Since both the number of diagonalizations and the number of cells would be increased by a factor of roughly  $5^3 = 125$ , we extended the linear approximation to include full quadratic terms. Volume integration of constant energy contours in the quadratic approximation is difficult, however, and involves the evaluation of many sets of incomplete elliptic integrals (as many as one has histogram boxes and cells). The procedure would be slow, difficult to program, and not easily extended to those cases where we want to consider  $\mathbf{k}$ -dependent matrix elements inside the integrand as well as singular energy delta functions.

The QUAD method may or may not be easier to use in more complicated calculations than the density of states in simple systems. Recently,

Saravia and Brust<sup>3</sup> successfully extended the linear-approximation method to evaluate matrix elements for the energy dependence of the long-wavelength dielectric function  $\epsilon_2(\omega)$  of diamond. We believe that the linear-approximation method cannot be successfully applied to materials with more convoluted  $\underline{k}$  dependence of electronic properties.

In this report we present two sets of Fortran programs used to treat fcc transition-metal bands. The physical and mathematical approximations involved in these programs have been presented in two companion articles.<sup>4,5</sup> The extension of this treatment to other crystal structures and materials should be straightforward.

## II. TREATMENT OF THE BRILLOUIN ZONE

The Hamiltonian of a cubic material has eigenvalues which are degenerate under the 48 operations of the cubic point group. The full BZ may therefore be factorized into 48 irreducible BZ wedges, each containing exactly the same eigenvalue information. Although seemingly advantageous, a direct factorization of the full zone into one irreducible wedge can involve several problems. First, if Monte Carlo sampling methods are used to find the density of states (as below), each pseudorandom point must be tested to see if it is inside the wedge. Such testing requires several Fortran IF statements, which are slow compared to arithmetic statements.

A second possible problem induced by the wedge planes occurs when a regular or an irregular interpolation net is constructed over one cell. In an fcc structure, a simple cubic net has points falling inside and outside of the irreducible wedge planes, forming various gap or defect regions. Although a more complicated interpolation net, formed by shrinking the reciprocal-lattice basis vectors, will not overlap the wedge planes, its use requires a nonorthogonal basis set to derive the expansion coefficients needed for interpolation. Use of nonorthogonal bases can either raise the fractional error of the expansion coefficients to unacceptable levels, or leave some of the expansion coefficients indeterminate. For example, the simple cubic net discussed below cannot be used to find expansion coefficients to full third order.

Our treatment of the BZ avoids both of these problems at the outset; nevertheless, it does not need to diagonalize the Hamiltonian at any point outside of a single 1/48th of the BZ.

The QUAD technique treats the fcc BZ in the following manner:

1. The planes forming the square faces are extended so that a new BZ, exactly twice the volume of the old, is formed. The new BZ is a cube. Note that all of the 14 Bravais lattices can, by suitable patching-on of

degenerate pieces, be formed into, first, an orthorhombic system and then, by a scale transformation, into a cubic system. Thus our treatment, here restricted to the fcc system, can be applied to other systems as well.

2. Only points in the positive octant are considered. The volume of the resulting cube is 1/4 of the volume of a single BZ.

3. The 1/2 BZ distance ( $\Gamma$ -X) is divided by three sets of (MESH-1) parallel equispaced planes. The working volume is thus divided into (MESH\*\*3) little cubes, called cells.

4. Each cell is further subdivided by three bisecting planes, forming 27 points at the corners, midedges, midfaces, and midpoint of one cell. Eigenvalues at the 27 points are used to find, by a least-squares procedure, a set of ten expansion coefficients to full quadratic order for each of the nine bands, derived from the model Hamiltonian appropriate to noble and transition metals.<sup>4</sup>

5. Each of the various cells is labeled by an index. Except for such labeling, all cells are treated by the program as logically equivalent. The sampling range for  $\underline{k}$  is remapped for each cell so that it has components which are in the range -1 to +1. This remapping procedure greatly facilitates debugging. Since each cell is treated the same, it is only necessary to check in detail that one cell is being properly processed. In addition, accuracy in the matrix inversion procedure is increased, reducing the error in the expansion coefficients from about 1% to about 0.001% (both calculations were done using double precision arithmetic). Also, with the remapping, the  $10 \times 10$  coefficient matrix inversion need be performed only once. In the present program the matrix inversion is done only once per run, reducing total execution time by a factor of about two.

6. Each of the subdivided points, which fall on a regular net of length  $(\Gamma$ -X)/(2\*MESH), is mapped back onto one particular irreducible wedge. Thus the necessary energy eigenvalues at these points in 1/48th of the BZ are found only once, are stored in the matrix EMESH, and are drawn out as needed by the cells in the expanded cubic BZ segment. No more diagonalizations need be performed using this method than in other techniques, and the wedge planes need not be considered at all.

7. Monte Carlo sampling points are found throughout one cell, the nine energy eigenvalues appropriate to each point are found from the expansion coefficients CEP, and the energies are sorted into the HIST matrix.

The QUAD technique may be easily extended to include  $\underline{k}$ -dependent scalar, vector, or tensor quantities. When  $\underline{k}$ -dependent elements are important, two QUAD expansions are used: one, as above, for the energy eigenvalues, and another for the eigenvectors. From the eigenvectors, the

basis functions, and the various form factors, any  $\underline{k}$ -dependent quantity may be calculated. Since the experimental accuracies are usually of the order of a few percent, about 10,000 Monte Carlo points per histogram box need be calculated to match these accuracies. Modern computers such as the Argonne 360/50/75 can calculate this many points in a few seconds, typically 500 to 1000 per second. Only 1/2 to 1 hour of machine time is needed to calculate a fine histogram of the most complicated experimental quantities. Once the basic band-structure parameters and various form factors are known, full band-structure details can be routinely included in the most difficult calculations using the QUAD technique.

### III. DESCRIPTION OF THE PROGRAM

The program presented here forms a fine histogram representation of the density of electron states of relativistic fcc transition metals. As we have suggested above, such a calculation may be viewed as a "process" in which all matrix elements are strictly equal to one or zero. The extension of the QUAD scheme techniques for either manipulating core space or improving running speed to problems where matrix elements have a strong  $\underline{k}$  dependence is straightforward. The program was written for the version 13, release 11 of Fortran IV(H) compiler of the IBM 360/50/75 system as the Computation Center of the Applied Mathematics Division at Argonne National Laboratory. This system has available about 250,000 single-precision words of fast core. The program now uses about 40,000 single-precision words, and 25,000 double-precision words of core. The need for core space could be substantially lowered by simply reducing the dimension of the storage matrices EMESH, IP, and HIST. Moreover, HOST, the integrated density of states, could be eliminated entirely.

#### A. Main Program

The main routine is used for manipulative purposes only, i.e., to call other routines, allocate storage space, transfer information, or call input-output. Statement cards are included at key points in the programs. Labeled common blocks are used to transmit information to the subroutines. The common blocks have been factored according to operational function.

#### B. Subroutines

The subroutines, with a brief description of each, are listed below in alphabetical order. Names containing an asterisk are used only by the combined interpolation scheme.

##### ALL

The ALL routine finds the energy eigenvalues for the nine bands of a fcc transition metal and loads the eigenvalues sequentially into EMESH,

the energy mesh. As discussed above, two nested meshes are used. An outer mesh labels the various cells of length  $\left(\frac{2\pi}{a} * \frac{1}{MESH}\right)$ ; an inner mesh of length  $\left(\frac{2\pi}{a} * \frac{1}{2*MESH}\right)$  is the regular eigenvalue net where a is the lattice constant. The inner mesh is assumed to be fcc and is generated in the first part of the ALL routine. The mesh points are given in integer arithmetic, and the integer vector [111] is added to them. For example, the point  $\Gamma$  becomes [111] and, assuming MESH = 4, the point X becomes [911]. The matrices IP and IK are logical inverses of each other and contain essentially the same information. The matrix IP yields the index point of a triplet set of numbers, whereas the matrix IK yields the k-space coordinates of a given index number. Use of both of these matrices allows immediate retrieval of information without searching arrays. Appreciable computation time and core space are saved, and the procedure can be applied to many other computational problems.

### CELL

The CELL routine returns the ten possible quadratic products of the vector BK. The first nine products (the harmonics s, p, and d) are scaled to agree with the Kubic Harmonics used in the interpolation scheme. This rescaling was done arbitrarily and represents an historical rather than a necessary choice.

### \*EIGEN B

The EIGEN B routine finds the highest NEV eigenvalues and eigenvectors of the matrix A to an accuracy of ACC. The secular equation is of order NSUB. If M = 0 only eigenvalues are found. These are stored in the array VALU, starting with the highest. EIGEN B uses the tridiagonal procedure of Givens and Householder, and was originally written by Burton Garbow of the Applied Mathematics Division at Argonne. It is considerably faster than the Jacobi technique.

### \*FILL

The FILL routine forms the spin-orbit matrix elements of the spin-up, spin-down d-functions. The original  $9 \times 9$  real secular equation of the interpolation scheme is extended to an  $18 \times 18$  complex matrix to include spin. The Hermitian Hamiltonian has been doubled in size to make a real matrix by the usual procedure:  $(A + iB) \rightarrow \begin{pmatrix} A & B \\ -B & A \end{pmatrix}$ , where A and B are block matrices. The spin-orbit coupling parameter is EP, and we assume a spin-orbit interaction of the form

$$H_{SO} = \frac{EP}{2} \vec{\sigma} \cdot \vec{L},$$

where  $\vec{L}$  acts on d states alone.

\*HELD

The HELD routine generates the Hamiltonian of the  $5 \times 5$  tight-binding fcc, d-d interaction block, and includes the effects of second-nearest neighbors. HELD is the result of straight coding of Table II of Slater and Koster.<sup>6</sup>

\*HSOC

The HSOC routine generates two blocks of the model Hamiltonian. The first is the diagonal  $4 \times 4$  block of the Hamiltonian between plane waves. The second is the off-diagonal  $5 \times 4$  block of the Hamiltonian between tight-binding d states and plane waves. The five d basis functions have the same order as in the HELD routine. The four plane waves (OPW's) are the lowest four which are degenerate in the point W in the BZ.

\*MATMPY

The MATMPY routine is a matrix multiplier used by EIGEN B.

MDINV

The MDINV routine calculates the inverse of the double-precision matrix A up to order N. A technique which pivots the matrix on the largest elements is used. The inverse is returned in the matrix A. The original A is destroyed.

\*RECIP

The RECIP routine finds that point W, of three possible W points in the positive truncated octant, which is closest to the given point BK. The components of the lowest four degenerate plane waves are stored in the matrix AK and form the plane-wave basis set used in HSOC. The effect of RECIP is cancelled out in the present program by the restrictions imposed by ALL. RECIP becomes important in other cases--particularly when the effects of external magnetic fields are introduced into the Hamiltonian or when eigenvectors rather than eigenvalues are desired. The RECIP routine was kindly lent to us by E. I. Zornberg,<sup>7</sup> who has made several modifications of the original scheme.<sup>4</sup>

RANF

The RANF routine generates a set of pseudorandom numbers uniformly distributed between 0.0 and 1.0. The numbers are generated by a technique using the multiplicative congruence

$$X_{n+1} = X_n(2^{16} + 1) \text{ modulo } 2^{31}.$$

This technique assumes an integer word width of 31 binary bits and left truncation of multiplicative overflow. The RANF function was written by Nancy W. Clark of the Applied Mathematics Division at Argonne. Modification to other computer systems should be straightforward and would be based on the integer word width of the system.

Several tests have been made of the accuracy of this pseudorandom generator. Three Chi square tests of the one-dimensional distribution of samples of 10,240 numbers with 1023 degrees of freedom gave the results 977, 1012, and 926. The 95% confidence interval for this test is (936, 1113).

Three Chi square tests of the three-dimensional distribution of triples of numbers from these same samples done with 511 degrees of freedom gave the results 485, 515, and 496. The 95% confidence interval for this test is (450, 577).

For our simple needs this routine considerably overkills the problem of generating pseudorandom numbers. However, other, less-precise techniques would not result in any substantial savings of computational time.

#### \*RF and \*RG

The RF and RG routines generate the two form factors  $\langle k|d \rangle$  and  $\langle k|H|d \rangle$ , respectively, used in the interpolation scheme<sup>4</sup> subroutine HSOC. The form of both functions is basically  $j_2(x)$ , the ordinary spherical Bessel function of order 2.

#### \*SET UP

The SET UP routine reads parameters into the various common block storage locations. The parameters provided in Table I are appropriate

TABLE I. The Data for the QUAD Scheme  
Programs for the Bands of Platinum

Parameter	Value	Parameter	Value
MESH	00010	P12	1.7928
ND	00060	P13	0.3700
HEAD	THE PLATINUM PARAMETERS	P14	0.2683
CNORM	0.9830	P15	0.0241
P1	-1.0689	P16	0.6545
P2	0.0027	P17	0.9745
P3	0.0033	P18	0.0010
P4	0.6768	P19	-0.0023
P5	-0.4961	P20	-0.0054
P6	-0.4845	P21	-0.0012
P7	-0.0636	MSIZE	00036
P8	+0.0254	NT	01000
P9	0.0193	TOP	0.8000
P10	-0.0003	BOT	0.2000
P11	1.9493		

to fcc platinum. The platinum energy bands have been fit to the accurate and extensive dHvA experiments of Ketterson and Windmiller.<sup>8</sup> Note that several of the parameters  $P_i$  have nonband structure functions, e.g.,  $P_{16}$  is the Fermi energy. The parameter  $P_{17}$  uniformly widens or narrows the d-band complex.

### SNORT

The SNORT routine finds the nine expansion coefficients CEP for the cell whose lowest corner point is given by IVEC. The element N69 in common block /HOT DOG/ is used as a test location to restrict matrix inversion of TR by MDINV to the first entry of SNORT. Every cell point used in SNORT has components restricted to  $\pm 1$  or 0.

### \*SYM

The SYM routine is used by the interpolation scheme to symmetrize the Hamiltonian calculated in HELD and HSOC.

### \*ZEP

The ZEP routine calls the various routines used by the combined interpolation scheme to give the nine energy eigenvalues appropriate to the point CK in the BZ. If the QUAD technique were used with a different band structure method, the ZEP routine could be replaced by the appropriate equivalent. Thus, the generalization of the present method to APW, KKR, and OPW band structure methods should be immediate.

## IV. FORTRAN LISTINGS

The Fortran listings of the main program and the subroutines follow:

```
C*****C
C      ..PROGRAM QUAD....
C*****C
C      J360--PERFORM MONTE CARLO INTEGRALS
C      SO AS TO OBTAIN DENSITY OF STATES
C      FOR FCC TRANSITION METALS USING
C      THE COMBINED INTERPOLATION SCHEME
C
C.....F.M.MUELLER, J.W.GARLAND, M.H.COHEN, AND K.H.BENNEMANN.....
C      JUNE 1968
C      MODIFIED BY S.G.DAS AUGUST 1969
C
C*****C
C      ND IS THE TOTAL NUMBER OF HISTOGRAM STATES
C      NT IS THE NUMBER OF POINTS PER BOX
C      TOP IS THE HIGHEST ENERGY VALUE
C      BOT IS THE LOWEST ENERGY VALUE
C*****C
C      FORMAT(F9.6)
1
```

```

4 FORMAT(15)
5 FORMAT(1H1,5X,'THE TOTAL NUMBER OF POINTS IS ',I5,/)
6 FORMAT(5X,'THE RANGE IS FROM ',F10.4,": TO ",F10.4,/)
7 FORMAT(1H1,/,10X,'THE DENSITY OF TOTAL ELECTRON STATES',/)
8 FORMAT(1H1,/,11X,'CELL',9X,'ENERGY',13X,'DENSITY OF STATES IN RY
1 AND EV',9X,'TOTAL STATES',/)
11 FORMAT(5X,215,5X,F10.5,10X,E15.5,5X,E15.5,5X,E15.5)
12 FORMAT(15,I5,4E15.4)
169 0 FORMAT(1H1,10X,'THE TOTAL NUMBER OF POINTS USED HERE IS ',I15,/
1//)
500 FORMAT(10X,3I15,/)
502 FORMAT(10X,9F10.4,/)
505 FORMAT(1H1,10X,'THE TOTAL NUMBER OF K SPACE POINTS IS ',I10,/) 00002250
2121 FORMAT(3F10.4) 00002300
2122 FORMAT(3E15.8)
C ****
C
C ****
C
C DIMENSION IVEC(3), CEP(9,10), DK(3), C(10)
DIMENSIONS M1(9), DMI1(9), SMI(9), DMI2(9), DMIT1(9), DMIT2(9)
DIMENS(ONDENS(2000), Y5(680)
DIMENSION DME21(9,300), DE21(9,10), SME1(9,300), SEP1(9,10), SME2(9,
8300), SEP2(9,10), DME31(9,300), DE31(9,10), DME32(9,300), DE32(9,
4300), DME22(9,300), DE22(9,10)
DIMENSIONS (340), Y1(680), Y2(680), Y3(680), Y4(340), X1(680)
DIMENSION EMESH(9,2000), IP(25,25,25), IK(3,2000)
1,CK(3),EG(9)
DIMENSION ADV(3)
C
REAL*8 HIST(13,500),HOST(13,500),HLOW(I3),SUM,SNUM,SC
C
COMMON/GURU/DME21,DE21,SME1,SEP1,SME2,SEP2,DME31,DE31,DME32,DE32,
7DME22,DE22,SMI,SMII,DMI1,DMI2,DMIT1,DMIT2 00003350
COMMON /DEBUG/ HIST,HOST
COMMON/HOTDOG/N69
COMMON/BLOK/EMESH,MESH,IK,NTOT,MC,IP
C
C ****
C
C SET CONSTANTS
N69 = 0
SQ3=SQRT(3.)
DO 111 I = 1,13
HLow(I) = 0.000
111 CONTINUE
C
C ****
C
C SET THE MATRIX OF MESH POINTS AND THE NUMBER OF CALLS
READ (5,4) MESH
READ (5,4) ND
XND=ND
C
C SET UP DATA FROM CARDS
CALL SET UP
C
PARAMETER DEFINITION COMPLETED
C
C ****
C
C FILL UP LATTICE POINTS
CALL ALL
C
C ****
C
WRITE (6,505) NTOT
DO 102 MC = 1,NTOT

```

```

      WRITE(6,500) (IK(I,NC),I = 1,3)
      WRITE(6,502) (EMESH (J6,NC),J6 = 1,9)
      WRITE(6,502) (DME21 (J6,NC),J6 = 1,9)
      WRITE(6,502) (DME22 (J6,NC),J6 = 1,9)
      WRITE(6,502) (DME31 (J6,NC),J6 = 1,9)
      WRITE(6,502) (DME32 (J6,NC),J6 = 1,9)
      WRITE(6,502) (SME1 (J6,NC),J6 = 1,9)
      WRITE(6,502) (SME2 (J6,NC),J6 = 1,9)
102   CONTINUE
C
C***** ****
C
C      SET UP SCALE
      IMESH=2*MESH+1
      XL=IMESH-1
      SCALE=8./XL
C
C      CLEAR HISTOGRAMS AND SET UP PARAMETERS
      DO62I=1,ND
      DO 62 J = 1,13
      HIST(J,I) = 0.000
      HOST(J,I) = 0.000
62   CONTINUE
C
C***** ****
C
      READ(5,4)NT
      READ(5,1)TOP,BOT
      DEL=(TOP-BOT)/ND
      NW = MESH*MESH*MESH*NT
      WRITE(6,169) NW
      WRITE(6,5)NT
      WRITE(6,6)BOT,TOP
C
C***** ****
C
      SNUM=0.
      MD=2*ND
C
C***** ****
C      BEGIN LOOP ON INDIVIDUAL LATTICE BOXES
C***** ****
C
      DO100IIX=1,MESH
      DO100IY=1,MESH
      DO100IZ=1,MESH
C
      IVEC(1)=IX
      IVEC(2)=IY
      IVEC(3)=IZ
C
      SET UP COEFFICIENTS
C
C***** ****
      CALLSNORT(IVEC,CEP)
C***** ****
C
      RUN ON INNER LOOP USING EXPANSION COEFFICIENTS
C
      DO50NC=1,NT
      SCALE UP VECTORS
      DO63I=1,3
      DK(I) = RANF(-1)*2-1
63   CONTINUE
      READY TO ADD AND CALL CELL
      CALLCELL(DK,C)
      LOAD IT IN 9 LEVELS
C

```

```

DO 40 JG=1,9
C
      ER=0.
      PD1=0.0
      PD2=0.0
      PD31=0.0
      PD32=0.0
      PS1=0.0
      PS2=0.0
C
      DO 30 IG=1,10
C
      ER=ER+C(IG)*CEP(JG,IG)
      PS1=PS1+C(IG)*SEP1(JG,IG)
      PS2=PS2+C(IG)*SEP2(JG,IG)
      PD1=PD1+C(IG)*DE21(JG,IG)
      PD2=PD2+C(IG)*DE22(JG,IG)
      PD31=PD31+C(IG)*DE31(JG,IG)
      PD32=PD32+C(IG)*DE32(JG,IG)
C*****CONTINUE*****
30    CONTINUE
C*****CONTINUE*****
C
      IF(ER.LT.BOT)GOTO85
      IF(ER.GT.TOP)GOTO80
      NSET=(ER-BOT)/DEL+1-
C
C*****CONTINUE*****
C      COUNT TWO ELECTRONS
C*****CONTINUE*****
C
      HIST(13,NSET)=HIST(13,NSET)+2.0D0*(PS1*PS2)
C
80    CONTINUE
      GO TO 86
85    CONTINUE
C
      HLOW(JG) = HLOW(JG)+2.0D0
      HLOW(10) = HLOW(10)+2.0D0
      HLOW(11) = HLOW(11)+2.0D0*(PD1+PD2)
      HLOW(12) = HLOW(12)+2.0D0*(PD31+PD32)
      HLOW(13) = HLOW(13)+2.0D0*(PS1+PS2)
C
86    CONTINUE
C
      SNUM=SNUM+2.0D0
C
40    CONTINUE
50    CONTINUE
100   CONTINUE
C*****CONTINUE*****
C      RENORMALIZE TO PROPER NUMBER OF ELECTRONS
C*****CONTINUE*****
C
      SC = 18.0D0/SNUM
C
      DO 110 J = 1,13
      SUM = HLOW(J)*SC
      DO 110 I=1,ND
      HIST(J,I)=HIST(J,I)*SC
      SUM=SUM+HIST(J,I)
      HOST(J,I) = SUM
C      CONVERT TO DENSITY
      HIST(J,I) = HIST(J,I)/DEL
110   CONTINUE
C

```

```

C*****SET UP THE OUTPUT*****
C*****DO 300 K = 1,13
      DO 300 K = 1,13
      WRITE(6,7)
      NG=ND/55
      DO200I=1,NG
      WRITE(6,10)
      NTOP=I*55
      NBOT=NTOP-54
      DO200J=NBOT,NTOP
      ENG= J-1
      ENG=ENG+DEL+BOT
      HS=HIST(K,J)/13.59
      WRITE(6,11)K,J,ENG,HIST(K,J),HS,HOST(K,J)
C     PUNCH H12,K,J,ENG,HIST(K,J),HS,HOST(K,J)
200   CONTINUE
300   CONTINUE
C
C*****STOP
      STOP
      END

```

#### SUBROUTINE ALL

```

C*****J353-FILLS ALL POSSIBLE LATTIC POINTS
C     NOTE THAT IK IS ONE HIGHER THAN ZERO
C.....F.M.MUELLER,J.W.GARLAND,M.H.COHEN,AND K-H.BENNEMANN.....
C     JUNE 1968
C     MODIFIED BY S.G.DAS AUGUST 1969
C*****MESH IS NUMBER OF CELLS FROM
C     GAMMA TO X
C     IK SORTS VECTORS,EMESH ARE EIGENVALUES
C     NTOT IS THE NUMBER OF DOUBLE CELL POINTS
C*****DIMENSION DME21(9,300),DE21(9,10),SME1(9,300),SEP1(9,10),SME2(9,
8300),SEP2(9,10),DME31(9,300),DE31(9,10),DME32(9,300),DE32(9,
4300),DME22(9,300),DE22(9,10)
DIMENSIONS M11(9),DM11(9),SM1(9),DM12(9),DMIT1(9),DMIT2(9)
DIMENSION EMESH(9,2000),IP(25,25,25),IK(3,2000)
DIMENSION BK(3),EG(9)
C
COMMON/GURU/DME21,DE21,SME1,SEP1,SME2,SEP2,DME31,DE31,DME32,DE32,
7DME22,DE22,SM1,SM11,DM11,DM12,DMIT1,DMIT2
COMMON/BLOK/EMESH,MESH,IK,NTOT,MC,IP
C
C*****DEFINE AND FILL-UP EMESH AND IK
C*****C
C     FOLLOWING ASSUMES F-C-C. LATTICE
C
IMESH=2*MESH+1
C     IK FIRST
INDEX=0
C
DO50IY=1,IMESH
DO50IX=1,IY
DO50IZ=1,IX
C
PUT IN L PLANE
IT=IX+IY+IZ-3

```

```

IF(MESH*3.LT.IT) GOTO50
C
INDEX=INDEX+1
I=INDEX
IP(IX,IY,IZ)=INDEX
IK(1,I)=IX
IK(2,I)=IY
IK(3,I)=IZ
C
50 CONTINUE
C
NTOT=INDEX
C
C*****IK NOW DONE*****
C
C      SET SCALE OF ZONE
C
XL=IMESH-1
SCALE=8./XL
C
C      LOAD EMESH
C
DO100I=1,NTOT
C
C      RESET VECTOR SO THAT GAMMA=ZERO
C
DO60J=1,3
BK(J)=(IK(J,I)-1)*SCALE
60 CONTINUE
C
C      CALL MATRIX AND EIGENVALUE ROUTINES
C
CALL ZEP(BK,EG)
C
C*****FILE INTO EMESH*****
C
C      DO70 J=1,9
C
EMESH(J,I) = EG(J)
SME1(J,I)=SM1(J)
SME2(J,I)=SM2(J)
DME31(J,I)=DM1T1(J)
DME32(J,I)=DM1T2(J)
DME21(J,I)=DM1L(J)
DME22(J,I)=DM1R(J)
C
70 CONTINUE
C
100 CONTINUE
C
C*****MESH LOADED*****
C
RETURN
END

SUBROUTINE CELL(BK,C)
C
C      J351-FOR USE IN DENSITY OF STATES
C      NOTE THAT A DIFFERENT SET OF EXPANSION FUNCTIONS

```

```
C      CAN BE USED BY MODIFYING THIS ROUTINE
C.....F.M.MUELLER, J.W.GARLAND, M.H.COHEN, AND K.H.BENNEMANN.....
C      JUNE 1968
C*****
C      DIMENSIONBK(3),C(10)
C
C*****SET UP VALUE AT THIS POINT OF KSPACE
C*****
C
C      SQ3=SQRT(3.)
C
C      C(1)=1.
C      C(2)=BK(1)
C      C(3)=BK(2)
C      C(4)=BK(3)
C      C(5)=BK(1)*BK(2)
C      C(6)=BK(1)*BK(3)
C      C(7)=BK(2)*BK(3)
C      C(8)=(BK(1)*BK(1)-BK(2)*BK(2))*5
C      C(9)=.5*(3.*BK(3)*BK(3)-1.)/SQ3
C      C(10)=BK(1)*BK(1)+BK(2)*BK(2)+BK(3)*BK(3)
C
C*****
C      RETURN
C
C      END
```

```
FUNCTION CONVO(X)
C
C*****THIS FUNCTION IS A DUMMY
C      OTHER ENTRY TO CONVERT FROM FORTRAN 2 TO FORTRAN 4
C*****
C
C      F.M.MUELLER MARCH 1966
C
C*****
C
C      CONVO = 1
C      RETURN
C
C      ENTRY ABSF(X)
C      ABSF = ABS(X)
C      RETURN
C
C      ENTRY COSF(X)
C      COSF = COS(X)
C      RETURN
C
C      ENTRY SINF(X)
C      SINF = SIN(X)
C      RETURN
C
C      ENTRY SQRTE(X)
C      SQRTE = SQRT(X)
C      RETURN
C
C      ENTRY SIGNF(X,Y)
C      SIGNF = SIGN(X,Y)
C      RETURN
C
C*****END
```

```

SUBROUTINE EIGENB(NSUB,NEV,M,ACC)
C **** MODIFICATION OF ANLF202 ****
C ****
C      DIMENSION A( 36, 36),B( 36, 36),VALU(36),VALUL(36),DIAG(36),SUPERD( 36), 3
C      1 Q(36),S(36),C(36),D(36),IND(36),U(36) 3
C      COMMON A,VALU 4
C      2 N=NSUB 5
C
C      CALCULATE NORM OF MATRIX 5
C
C      3 ANORM2=0.0 6
C      4 DO 6 I=1,N 7
C      5 DO 6 J=1,N 8
C      IF (ABS(A(I,J)).LT.1E-12)A(I,J)=0.0
C      6 ANORM2=ANORM2+A(I,J)**2 9
C
C      GENERATE IDENTITY MATRIX 9
C      7 ANORM=SQRT (ANORM2) 10
C      9 IF (M)10,45,10 11
C      10 DO 40 I=1,N 12
C      12 DO 40 J=1,N 13
C      20 IF ((I-J)35.25,35 14
C      25 B(I,J)=1.0 15
C      30 GO TO 40 16
C      35 B(I,J)=0.0 17
C      40 CONTINUE 18
C
C      PERFORM ROTATIONS TO REDUCE MATRIX TO JACOBI FORM 18
C
C      DIMENSION WVEC(36),PVEC(36),QVEC(36) 19
C      45 CONTINUE 20
C      50 NN=N-2 21
C      52 IF (NN)890,170,55 22
C      55 DO 160 I=1,NN 23
C      58 II=I+1 24
C      60 IZ=I+2 25
C      65 SUM=0.0 26
C      68 DO 70 J=IZ,N 27
C      70 SUM=SUM+A(I,J)**2 28
C      72 IF (SUM)75,160, 75 29
C      75 SUM=SQRT (SUM+A(I,II)**2) 30
C      80 WVEC(II)=SQRT ( 1.0+ABS (A(I,II))/SUM) 31
C      82 DIV=SIGN ( SUM+WVEC(II),A(I,II)) 32
C      83 DO 85 J=IZ,N 33
C      85 WVEC(J)=A(I,J)/DIV 34
C      88 NI=N-I 35
C      90 CALL MATMPY (NI,NI,WVEC(II),A(II,II),PVEC(II)) 36
C      95 CALL MATMPY (1,NI,WVEC(II),PVEC(II),SCALAR) 37
C      98 SCALAR=0.5*SCALAR 38
C      100 DO 105 J=II,N 39
C      105 QVEC(J)=PVEC(J)-SCALAR*WVEC(J) 40
C      112 A(I,II)=-SIGN (SUM,A(I,II)) 41
C      114 DO 120 K=II,N 42
C      116 DO 120 L=II,K 43
C      118 A(K,L)=A(K,L)-(WVEC(K)*QVEC(L)+WVEC(L)*QVEC(K)) 44
C      IF (ABS(A(K,L)).LT.1E-12)A(K,L)=0.0
C      120 A(L,K)=A(K,L) 45
C      125 IF (M)130,160,130 46
C      130 CALL MATMPY (N,NI,WVEC(II),B(II,1),QVEC) 47
C      140 DO 150 K=II,N 48
C      145 DO 150 L=1,N 49
C      150 B(K,L)=B(K,L)-WVEC(K)*QVEC(L) 50
C      160 CONTINUE 51
C
C      MOVE JACOBI FORM ELEMENTS AND INITIALIZE EIGENVALUE BOUNDS 51

```

```

C
170 DO 200 I=1,N
  VALU(I)=ANORM
  VALL(I)=-1.*ANORM
200  DIAG(I)=A(I,I)
210 DO 230 I=2,N
220 SUPERD(I-1)=A(I-1,I)
230 Q(I-1)=(SUPERD(I-1))**2
C
C      DETERMINE SIGNS OF PRINCIPAL MINORS
C
235 TAU=0.0
240 I=1
260 MATCH=0
270 T2=0.0
275 T1=1.0
277 DO 450 J=1,N
280 P=DIAG(J)-TAU
290 IF(T2) 300, 330, 300
300 IF(T1) 310, 370, 310
310 T=P*T1-Q(J-1)*T2
320 GO TO 410
330 IF(T1) 335, 350, 350
335 T1=-1.0
340 T=-P
345 GO TO 410
350 T1=1.0
355 T=P
360 GO TO 410
370 IF(Q(J-1)) 380, 350, 380
380 IF(T2) 400, 390, 390
390 T=-1.0
395 GO TO 410
400 T=1.0
C
C      COUNT AGREEMENTS IN SIGN
C
410 IF(T1) 425, 420, 420
420 IF(T) 440, 430, 430
425 IF(T) 430, 440, 440
430 MATCH=MATCH+1
440 IF (ABS (T)-1.0E20) 445,445,442
442 T1=T1/T
443 T=1.0
445 T2=T1
450 T1=T
C
C      ESTABLISH TIGHTER BOUNDS ON EIGENVALUES
C
460 DO 530 K=1,N
465 IF (K-MATCH) 470, 470, 520
470 IF(TAU-VALL(K)) 530, 530, 480
480 VALL(K)=TAU
490 GO TO 530
520 IF(TAU-VALU(K)) 525, 530, 530
525 VALU(K)=TAU
530 CONTINUE
540 IF ((VALU(I))-VALL(I)-ACC)570,570,550
550 IF((VALU(I))) 560, 580, 560
560 IF (ABS (VALU(I)/VALU(I)-1.0)-ACC)570,570,580
570 I=I+1
575 IF ((I-NEV)540,540,590
580 TAU=(VALL(I)+VALU(I))/2.0
585 GO TO 260
C
C      JACOBI EIGENVECTORS BY ROTATIONAL TRIANGULARIZATION
C
590 IF (M)593, 890, 593

```

```

593 CONTINUE          107
595 DO 610 I=1,N      108
600 DO 610 J=1,N      109
610 A(I,J)=0.0        110
615 DO 850 I=1,N      111
620 IF (I-1)625,625,621 112
621 IF (VALU(I-1)-VALU(I))-5.0E-7)730,730,622 113
622 IF (VALU(I-1))623,625,623 114
623 IF (ABS (VALU(I)/VALU(I-1)-1.0)-5.0E-7)730,730,625 115
625 COS=1.0           116
628 SIN=0.0           117
630 DO 700 J=1,N      118
635 IF (J-1)680,680,640 119
640 T=SQRT (T1**2+T2**2) 120
645 COS=T1/T           121
648 SIN=T2/T           122
650 S(J-1)=SIN         123
660 C(J-1)=COS         124
670 D(J-1)=T1*COS+T2*SIN 125
680 T1=(DIAG(J)-VALU(I))+COS-BETA*SIN 126
690 T2=SUPERD(J)       127
897 FORMAT(BE12.4)      128
700 BETA=SUPERD(J)*COS 129
710 D(N)=T1             130
720 DO 725 J=1,N        131
725 IND(J)=0            132
730 SMALLD=ANORM        133
735 DO 780 J=1,N        134
740 IF (IND(J-1))750,780,780 135
750 IF (ABS (SMALLD)-AHS (D(J)))780,780,760 136
760 SMALLD=D(J)         137
770 NN=J                138
780 CONTINUE            139
790 IND(NN)=1            140
800 PRODS=1.0            141
805 IF (NN-1)810,850,810 142
810 DO 840 K=2,NN        143
820 II=NN+1-K            144
830 A(II+1,I)=C(II)*PRODS 145
840 PRODS=PRODS*S(II)    146
850 A(1,I)=PRODS        146
C   FORM MATRIX PRODUCT OF ROTATION MATRIX WITH JACOBI VECTOR MATRIX 146
C   *****                                                               146
C   *****                                                               146
855 DO 885 J=1,N        147
860 DO 865 K=1,N        148
865 U(K)=A(K,J)
885 CALL MATMPY(N,N,U,B,A(1,J))
890 RETURN
END

```

## SUBROUTINE FILL

```

C***** ****
C   RETURNS THE SYMMETRIZED SPIN ORBIT MATRIX
C   F.M. MUELLER MARCH 1966
C***** ****
C   DIMENSION AD(12),AL(20,9),DIFER(20),BK(3)
C   DIMENSION SOB(36,36),EV(36),D(5)
C
C   COMMON SOB,EV,V0,V1,V2,V3,V4,ENGL,D,T,EIJ,R0,R1,EP,AD,AL,DIFER,BK
C***** ****
C   C=SQRTF(3.)
C
C   SET UP THE IMAGINARY PARTS

```

```

C
SOB(1,22) = 2.*EP
SOB(1,29) = -EP
SOB(2,21) = -EP
SOB(2,28) = EP
SOB(3,20) = EP
SOB(3,31) = -EP
SOB(3,32) = -C*EP
SOB(4,19) = -2.*EP
SOB(4,30) = EP
SOB(5,30) = C*EP
SOB(10,20) = -EP
SOB(10,31) = -2.*EP
SOB(12,22) = -EP
SOB(12,23) = -C*EP
SOB(12,29) = -EP
SOB(11,19) = EP
SOB(11,30) = EP
SOB(13,21) = EP
SOB(13,28) = 2.*EP
SOB(14,21) = C*EP
C
C*****SET UP THE REAL PARTS*****
C
C
SOB(1,12) = EP
SOB(3,10) = -EP
SOB(2,13) = -EP
SOB(2,14) = C*EP
SOB(4,11) = EP
SOB(5,11) = -C*EP
SOB(19,30) = EP
SOB(21,28) = -EP
SOB(20,31) = -EP
SOB(20,32) = C*EP
SOB(22,29) = EP
SOB(23,29) = -C*EP
C
DO 30 I = 1,9
DO 30 J = I,9
DO 30 K = 1,3
C
M = I+9*K
N = J+9*K
C
30 SOB(M,N) = SOB(I,J)
C
C*****CALL SYM(SOB,36)*****
C
C
RETURN
C
END

SUBROUTINE HELD
C
C*****SUBROUTINE TO GENERATE D PORTION OF HAMILTONIAN*****
C
C
THE ORDER OF THE BASIS STATES IS GIVEN BY.....
C
THE FIRST IS XY
C
THE SECOND IS XZ
C
THE THIRD IS YZ

```

```

C THE FOURTH IS (X*X-Y*Y)/2
C THE FIFTH IS (3*Z*Z-1)/(2*SQRT(3) )
C
C ****
C
C F.M.MUELLER MARCH 1966
C
C ****
C DIMENSION PREM(16),BK(3),D(5),AD(12),S(36,36),EV(36),EG(9)
C DIMENSION AL(20,9),DIFER(20)
C 1, CK(3)*EZ(9)
C COMMON S,E,V,V0,V1,V2,V3,V4,ENGL,D,T,EL,R0,R1,EP,A ,AL,DIFER,AK
C ,CNORM,MSIZE
C ****
C
C SET CONSTANTS
C
C PI = 3.14159265
C B = SQRTF(3.0)
C
C SET VARIABLES
C
C X = PI*AK(1)/8.
C Y = PI*AK(2)/8.
C Z = PI*AK(3)/8.
C
C SET COS,SIN FUNCTIONS
C
C CX = COSF(X)
C CY = COSF(Y)
C CZ = COSF(Z)
C SX = SINF(X)
C SY = SINF(Y)
C SZ = SINF(Z)
C C2X = COSF(2.*X)
C C2Y = COSF(2.*Y)
C C2Z = COSF(2.*Z)
C S2X = SINF(2.*X)
C S2Y = SINF(2.*Y)
C S2Z = SINF(2.*Z)
C
C ****
C GENERATION OF THE MATRIX ELEMENTS
C ****
C
C THE DIAGONAL TERMS
H(1,1) = A(1)+4.*A(2)*CX*CY+4.*A(3)*(CK* CZ+CY*CZ)
1+2.*A(9)*(C2X+C2Y)*2.*A(10)*C2Z
H(2,2) = A(1)+4.*A(2)*CX*CY+4.*A(3)*(CK* CY+CZ*CY)
1+2.*A(9)*(C2X+C2Z)*2.*A(10)*C2Y
H(3,3) = A(1)+4.*A(2)*CY*CZ+4.*A(3)*(CK* CY+CZ*CK)
1+2.*A(9)*(C2Y+C2Z)*2.*A(10)*C2X
H(4,4) = A(6)+3.*A(7)*(CX*CZ+CY*CZ)+A(8)*(4.*CX*CY+CX*CZ +CY*CZ)
1+(1.5)*A(11)*(C2X+C2Y)+(0.5)*A(12)*(C2X+C2Y+4.*C2Z)
H(5,5) = A(6)*A(7)*(4.*CX*CY+CX*CZ+CY*CZ)+3.*A(8)*(CX*CZ+CY*CZ)
1+(0.5)*A(11)*(C2X+C2Y+4.*C2Z)+(1.5)*A(12)*(C2X+C2Y)
C
C ****
C THE OFF-DIAGONAL TERMS
C
H(1,2) = -4.*A(4)*SY*SZ
H(1,3) = -4.*A(4)*SX*SZ
H(1,4) = 0.
H(1,5) = -4.*A(5)*SK*SY
H(2,3) = -4.*A(4)*SX*SY
H(2,4) = 2.*B*(5)*SX*SZ
H(2,5) = 2.*A(5)*SK*SZ

```

```

H(3,4) = -2.*B*A(5)*SY*SZ
H(3,5) = 2.*A(5)*SY*SZ
H(4,5) = B*A(7)*(CX*CZ-CY*CZ)-B*A(8)*(CX+CZ-CY*CZ)
1+(0.5)*B*(A(11)-A(12))*(CZY-CZX)
C
C ****
C
C      RETURN
C      END

SUBROUTINE HSOC
C
C ****
C      THIS ROUTINE CALCULATES THE OPW-D STATE INTERACTION BLOCKS
C      F.M.MUELLER MARCH 1966
C ****
C
C      DIMENSION GG(4),GF(4),GN(4),ES(5,4),AK(4,3)
C      DIMENSION S(36,36),EV(36),D(5),DIFER(20,BK(3)),AL(20,9),AD(12)
C      DIMENSION CK(3),ZK(4,3)
C
C      COMMONS, EV,V0,V1,V2,V3,V4,ENGL
C      COMMONS,T,EL,RB,RA,EP,AD,AL,DIFER,BK
C      COMMON/THIEF/AK
C ****
C
C      CALL RECIP(BK)
C
C      DO 2 I = 1,4
C
C      GO = 0.
C
C      DO 3 J = 1,3
C      CK(J) = AK(I,J)+BK(J)
C      ZK(I,J) = CK(J)
C 3    GO = CK(J)*CK(J) + GO
C      GI = SQRT (GO)
C
C      TEST=ABS(GI)
C      BEST=10.E-20
C
C      IF (TEST.GE.BEST) GO TO 94
C
C      GI=1.
C
94    CONTINUE
C
C      A = CK(1)/GI
C      B = CK(2)/GI
C      C = CK(3)/GI
C
C      ES(1,I) = A*B
C      ES(2,I) = A*C
C      ES(3,I) = B*C
C      ES(4,I) = .5*(A*A-B*B)
C      ES(5,I) = .5*SQRT (1./3.)*(3.*C*C-1.)
C
C      CALL RG(GI,Z,T,RB)
C
C      M = I+5
C
C      DO 50 J = 1,5
C      S(J,M) = ES(J,I)*Z
C
C      GG(I) = Z
C
C      CALL RF(GI,Z,EL,RA)

```

```

C
      GF(I) = Z
      GN(I) = SQRT (1.-GF(I)*GF(I)/3.)
C
      S(M,M) = GO*ENGL/48.+V0
C
C *****
C     DEFINE OFF DIAGONAL ELEMENTS
C *****
C
      S(6,7) = V1
      S(6,8) = V2
      S(6,9) = V1
      S(7,8) = V1
      S(7,9) = V2
      S(8,9) = V1
C
      DO 10 I = 6,9
C
      M = I-5
C
      DO 10 J = I,9
C
      N = J-5
C
      SUM = 0.
      SIM = 0.
C
      DO 11 K = 1,5
C
      SIM = SIM+ES(K,M)*ES(K,N)
C
      DO 11 L = 1,5
      A = S(K,L)
      11 SUM = SUM+ES(K,M)*ES(L,N)*A
C
      OS(I,J) = S(I,J)-SUM*GF(M)*GF(N)+SIM*(GN(M)*GG(M)*GF(N)
      1+GN(N)*GG(N)*GF(M))
C
      10   S(I,J) = S(I,J)/(GN(M)*GN(N))
C
C *****
C     RETURN
C *****
C
      END

```

```

SUBROUTINE MATMPY(M,N,X,A,Y)
C
C *****
C     THIS ROUTINE IS A HELPER TO EIGENB
C *****
C
      DIMENSION A(36,36),X(36),Y(36)
C
      10 DO 40 I=1,M
C
      20 Y(I)=0.0
C
      30 DO 40 J=1,N
C
      40 Y(I)=Y(I)+X(J)*A(J,I)
C
      50 RETURN
C *****
C
      END

```

## SUBROUTINE MDINV

```

C
C*****J115 DOUBLE PRECISION MATRIX INVERTER
C      NOTE THAT THIS MATRIX INVERTER CAN ONLY BE USED ON NEARLY DIAGONAL
C      MATRICIES
C      WE ARE NOT RESTRICTED TO TO SYMMETRIC MATRICIES
C*****
C      IMPLICIT REAL*8(A-E)
C*****
C
C      DIMENSION A( 35,35),IRL(35),ICL(35)
C      COMMON/OVER/A,N
C
C*****SET THE VARIOUS CALLS
C*****
C
M = N
M1 = M-1
IC = 0
C
DO 1 I = 1,M
IRL( I ) = 99
ICL(I) = I
1 CONTINUE
C
APIV = 0
BPIV = 1.0E20
ASUM = 0.
C
DO 2 I = 1,M
DO 2 J = 1,M
C
AT = DABS(A(I,J))
C
IF(AT.LT.1.E-30) AT = 0.
C
ASUM = ASUM+AT*AT
C
2 CONTINUE
C
IF(ASUM.LT.1.E-30) GO TO 5000
C
ASUM = DSQRT(ASUM)
ASUM = 1./ASUM
C
DO 3 I = 1,M
DO 3 J = 1,M
A(I,J) = A(I,J)*ASUM
3 CONTINUE
C
C*****FIND THE LARGEST ELEMENT
C*****
C
DO 10 I = 1,M
DO 11 J = 1,M
C
A1=A(I,J)
A2 = DABS(A1)
C
IF(A2.LE.APIV) GO TO 11
C
APIV = A2
C
IT = I

```

```
IU = J
C
11  CONTINUE
10  CONTINUE
C
C***** BEGIN THE ITERATION SET *****
C
C
24  CONTINUE
C
IF(APIV.EQ.0.) GO TO 14
IF(BPIV.LE.APIV) GO TO 15
C
BPIV = APIV
C
15  IR = IT
IS = IU
ITEMP = IRL(IR)
IRL(IR) = ICL(IS)
ICL(IS) = ITEMP-100
C
APIV = 0.
C
DIV = A(IR,IS)
A(IR,IS) = -1.
C
DO 17 J = 1,N
A(IR,J) = -A(IR,J)/DIV
C
DO 18 I = 1,M
C
IF(I.EQ.IR) GO TO 18
C
19  AIS = A(I,IS)
C
A(I,IS) = 0.
C
DO 20 J = 1,M
A(I,J) = A(I,J)+AIS*A(IR,J)
C
IF(J.GT.M) GO TO 20
IF(IRL(I).LE.50) GO TO 20
IF(ICL(J).LT.1) GO TO 20
C
AG = DABS(A(I,J))
C
IF(AG.LT.APIV) GO TO 20
APIV = AG
IT = I
IU = J
20  CONTINUE
C
18  CONTINUE
C
IC = IC+1
C
GO TO 24
C
C***** THE END *****
C
14  CONTINUE
C
DO 25 J = 1,M1
C
```

```
K = 0
KK = J+1
JSM = ICL(J)
C
DO 26 JJ = KK,M
C
IF(JSM.LE. ICL(JJ)) GO TO 26
C
K = JJ
JSM = ICL(JJ)
C
26 CONTINUE
C
IF(K.EQ.0) GO TO 25
C
DO 29 I = 1,M
C
ATP = A(I,J)
A(I,J) = A(I,K)
A(I,K) = ATP
C
29 CONTINUE
C
ITEMP = ICL(J)
ICL(J) = ICL(K)
ICL(K) = ITEMPI
C
25 CONTINUE
C
DO 30 I = 1,M1
C
K = 0
KK = I+1
ISM = IRL(I)
C
DO 31 II = KK,M
C
IF(ISM.LE. IRL(II)) GO TO 31
C
K = II
ISM = IRL(II)
C
31 CONTINUE
C
IF(K.EQ.0) GO TO 30
C
DO 34 J = 1,M
C
ATP = A(I,J)
A(I,J) = A(K,J)
A(K,J) = ATP
C
34 CONTINUE
C
ITEMP = IRL(I)
IRL(I) = IRL(K)
IRL(K) = ITEMPI
C
30 CONTINUE
C
C*****RESCALE THE MATRIX*****
C
DO 50 I = 1,M
DO 50 J = 1,M
C
A(I,J) = A(I,J)*ASUM
```

```
C
50      CONTINUE
C
C*****+
RETURN
C*****+
C
50 00  CONTINUE
C
      WRITE(6,100)
100  FORMAT(10X,'THESE MATRIX ELEMENTS ARE TOO SINGULAR',//)
C
C*****+
RETURN
C*****+
C
END

FUNCTION RANF(J)
C
C*****+
C      RANDOM NUMBER GENERATOR OF FORM X(I+1)=K(I)*(2**16+1) MOD 2**31
C      NOTE THAT THIS GENERATOR WORKS ON IBM 360 SYSTEM ONLY
C      MODIFICATION TO OTHER SYSTEMS IS STRAIGHT-FORWARD, HOWEVER
C      REPLACE 2**16 AND 2**31 BY RELATED WIDTH IN OTHER SYSTEM
C      WRITTEN BY NANCY CLARK OF AMD
C*****+
C
      EQUIVALENCE (X,IX)
      DATA IX/3125/
C
      2 IX=IX*65547
C
      IF(IX)5,6,6
C
      5 IX=IX+2147483647*1
C
      6 IF (J.GE.0) GOTO 8
C
      IF (IX.LT.8388608) GOTO 7
C
      YFL=IX/8
C
      RANF=YFL*.3725291E-8
C
C*****+
RETURN
C*****+
C
      7 YFL=IX
      RANF=YFL*.4656613E-9
C
C*****+
RETURN
C*****+
C
      8 RANF=X
C
C*****+
RETURN
C*****+
C
C*****+
ENTRY IRANF(J)
```

```
C*****
C      12 IX=IX*65547
C
C      IF(IX)15,16,16
C
C      15 IX=IX*2147 48 36 47*1
C      16 IRANF=IX
C
C*****
C      RETURN
C*****
C
C
C
C*****
C      ENTRY RANSET(J)
C*****
C
C      IX=J
C      RANSET=0.
C
C*****
C      RETURN
C*****
C
C
C
C*****
C      ENTRY RANGET(J)
C*****
C      J=IX
C      RANGET=0.
C
C*****
C      RETURN
C*****
C      END

SUBROUTINE RECIP(EK)
C*****
C      THIS MODIFICATION TO COMBINED INTERPOLATION SCHEME
C      SO THAT BK MAY BE IN ANY PART OF THE BZ
C      INTRODUCED BY E.I.ZORNBERG
C*****
C
C      DIMENSION AK(4,3),BK(3),EK(3),IL(3)
C      COMMON/THIEF/AK
C
C*****72    CONTINUE
C
C      DO 10 I=1,3
C      BK(I)=ABS(EK(I))
C
C      TES=BK(1)+BK(2)+BK(3)
C
C      IF(TES.LE.12.) GO TO 73
C
C      DO 13 I=1,3
C      EK(I)=EK(I)-SIGN(8., EK(I))
C
C      GO TO 72
```

```

C
C*****CONTINUE*****
73  CONTINUE
C*****CONTINUE*****
C
      DO 11 I=1,3
11    IL(I)=I
C
      DO 12 I=1,4
C
      DO 12 J=1,3
C
12    AK(I,J)=0.
C
      DO 70 I = 1,2
      DO 70 J = 2,3
C
      TEST = BK(I)-BK(J)
C
      IF(TEST.LE.0.) GO TO 70
C
      SAVE = BK(I)
      BK(I) = BK(J)
      BK(J) = SAVE
      KAVE=IL(I)
      IL(I)=IL(J)
      IL(J)=KAVE
C
70  CONTINUE
C
      BK(2) = BK(3)
      SAVE = BK(1)
      BK(1) = BK(2)
      BK(3) = SAVE
      KAVE=IL(1)
      IL(1)=IL(2)
      IL(2)=IL(3)
      IL(3)=KAVE
      I1=IL(1)
      I2=IL(2)
      I3=IL(3)
C
      IF(BK(I2).LE.8.) GO TO 71
C
      E2=SIGN(1.,EK(I2))
      EK(I2)=EK(I2)-E2 *16.
C
      GO TO 72
C
71  CONTINUE
C
      E2=SIGN(1.,EK(I2))
      AK(3,I2)=-E2*16.
      AK(2,I3)=-8.
      AK(4,I3)=8.
C
      DO 14 I=2,4,2
C
      AK(I,I1)=-SIGN(8.,EK(I1))
14    AK(I,I2)=-E2*8.
C
C*****RETURN*****
C*****END*****
      RETURN
C
      END

```

```

SUBROUTINE RF(GI,Z,T,RR)
C
C*****THIS ROUTINE CALCULATES THE OVERLAP FORM FACTOR
C      F.M.MUELLER MARCH 1966
C*****
C
31    A=GI*RR
C
IF(A-.8)3,3,32
C
3     Z = (3./A**3-1./A)*SIN (A)-3.*COS (A)/A**2
Z = T*ABS (Z)
C
GO TO 16
C
32    Z=0.
C
C*****
16    RETURN
C*****
C
END

SUBROUTINE RG(GI,Z,T,RR)
C
C*****THIS ROUTINE CALCULATES THE HYBRIDIZATION FORM FACTOR
C      F.M.MUELLER MARCH 1966
C*****
C
C = T
A=GI*RR
C
IF(A-.25) 20,20,10
C
10   IF(A-.5) 16,15,15
C
15   Z = 0.
C
C*****
RETURN
C*****
C
16   C = T*(5.1-A)/.85
20   Z = (3./A**3-1./A)*SIN (A)-3.*COS (A)/A**2
Z = Z*C
C
C*****
RETURN
C
END

SUBROUTINE SETUP
C
C*****J355 -- TO READ IN THE DATA FROM THE BACK DECK
C      AND SET UP ALL OF THE VARIOUS PARAMETERS
C      F.M.MUELLER MARCH 1966
C*****
C
100  FORMAT(15)
200  FORMAT(3E15.6)
300  FORMAT(F9.6)
350  FORMAT(10X,I5,10X,F10.4)
600  FORMAT(80A1)

```

```

700  FORMAT(1H1,10X,80A1,///)
800  FORMAT(//,10X,"D BAND PARAMETERS")
C
C*****+
C      DIMENSION HEAD(80)
DIMENSION AL(20,9),DIFER(20)
DIMENSION PREM(21),BK(3),D(5),AD(12),S(36,36),EV(36),EG(9)
COMMON S,E,V,V0,V1,V2,V3,V4,ENGL,D,T,EL,R0,R1,EP,AD,AL,DIFER,BK
1,CNORM,MSIZE
C
C*****+
C      READ(5,600) HEAD
WRITE(6,700) HEAD
C
C*****+
C      READ VALUES OF PARAMETERS
C*****+
C      READ (5,300) CNORM
READ (5,300) PREM
C      I = 0
C      WRITE(6,350) I,CNORM
WRITE(6,350)(I,PREM(I),I = 1,21)
C      READ (5,100) MSIZE
C
C*****+
C      SET UP PARAMETERS
C*****+
C      V0=PREM(1)
V1=PREM(2)
V2=PREM(3)
ENGL=PREM(4)
C
DO 20 I=1,4
I1=I+4
20 D(I)=PREM(I1)
D(5)=PREM(10)
C
T=PREM(11)
EL=PREM(12)
R0=PREM(13)
R1=PREM(14)
EP=PREM(15)
EF=PREM(16)
V3=0.
V4=0.
C
C*****+
C      VARY THE WHOLE WIDTH OF THE D BAND
C*****+
C      DO 169 JJ = 3,5
C
D(JJ) = D(JJ)*PREM(17)
C
169  CONTINUE
C
C*****+
C      SET UP D BAND PARAMETERS
C*****+

```

```

AD(1)=D(1)
AD(2)=-.25*(3.*D(3)+D(5))
AD(3)=.5*(D(4)+D(5))
AD(4)=-.5*(D(4)-D(5))
AD(5)=-.25*SQRT(3.)* (D(3)-D(5))
AD(6)=D(2)+D(1)
AD(7)=-.25*(D(3)+3.*D(5))
AD(8)=PREM(9)*PREM(17)
C      SET THE SECOND NEIGHBOR PARAMETERS
AD(9) = PREM(18)*PREM(17)
AD(10) = PREM(19)*PREM(17)
AD(11) = PREM(20)*PREM(17)
AD(12) = PREM(21)*PREM(17)
C
C      WRITE(6,800)
C      WRITE(6,200) AD
C
C*****RETURN*****
C*****END*****

```

```

SUBROUTINE SNORT(IVEC,CEP)
C
C*****J354-FINDS EXPANSION COEF'S TO SECOND ORDER
C      EMESH AND IK ALREADY LOADED
C.....F.M.MUELLER, J.W.GARLAND, M.H.COHEN, AND K.H.BENNEMANN.....
C      JUNE 1968
C      MODIFIED BY S.G.DAS AUGUST 1969
C*****IVEC IS LOWEST CORNER POINT IN
C      MESH-SIZED LATTICE
C
C
C      DIMENSION DME21(9,300),DE21(9,10),SME1(9,300),SEP1(9,10),SME2(9,
8300),SEP2(9,10),DME31(9,300),DE31(9,10),DME32(9,300),DE32(9,
4300),DME22(9,300),DE22(9,10)
DIMENSIONS MII(9),DMI1(9),SMI(9),DMII2(9),DMIT1(9),DMIT2(9)
DIMENSION EMESH(9,2000),IP(25,25,25),IK(3,20,2000)
DIMENSION IVEC(3),ISVEC(3),IZ(3),IR(3)
DIMENSION NEG(27),Z(3),G(10),IM(3,27)
DIMENSION CEP(9,10)
REAL U(10,27)
C
REAL*8 TR(35,35),R(27,10)
C
COMMON/OVER/TR,NGO
COMMON/BLOK/EMESH,IK,NTOT,MC,IP
COMMON/GURU/DME21,DE21,SME1,SEP1,SME2,SEP2,DME31,DE31,DME32,DE32,
7DME22,DE22,SMI,SMII,DMII,DMII2,DMIT1,DMIT2
COMMON/HOTDOG/N69
C
C*****NGO = 10
C
C      INITIATE THE CEP'S TO ZERO
C
DO 1 I = 1,9

```

```

      DO 1 J = 1,10
C
      CEP(I,J) = 0.
      SEP1(I,J)=0.0
      SEP2(I,J)=0.0
      DE21(I,J)=0.0
      DE22(I,J)=0.0
      DE32(I,J)=0.0
      DE31(I,J)=0.0
C
      1 CONTINUE
C
      DO4I=1,10
      DO3J=1,27
      R(J,I)=0.
C
      4 CONTINUE
C
C*****SET SCALE FOR DOUBLE SIZED MESH
C*****SET SCALE FOR DOUBLE SIZED MESH
C*****SET SCALE FOR DOUBLE SIZED MESH
C
      IMESH=2*MESH+1
      XL=IMESH-1
      SC AL E=8./XL
C
C*****PUT IVEC INTO DOUBLE LATTICE
C*****PUT IVEC INTO DOUBLE LATTICE
C*****PUT IVEC INTO DOUBLE LATTICE
C
      DO10I=1,3
      ISVEC(I) = 2*(IVEC(I)-1)
C
      10 CONTINUE
C
C*****SET UP TRIPLE-LOOP FOR 27 CELL POINTS
C*****SET UP TRIPLE-LOOP FOR 27 CELL POINTS
C*****SET UP TRIPLE-LOOP FOR 27 CELL POINTS
C
      NR=0
C
      DO20I=1,3
      DO20J=1,3
      DO20K=1,3
C
C*****IF WANT RESTRICTION TO FCC PUT HERE
C*****IF WANT RESTRICTION TO FCC PUT HERE
C*****IF WANT RESTRICTION TO FCC PUT HERE
C
      IZ(1) = I-2
      IZ(2) = J-2
      IZ(3) = K-2
C
      NR = NR+1
C
      DO15L=1,3
C
      IR(L) = IZ(L)
      IM(L,NR)=IR(L)
      Z(L)=IR(L)
C
      15 CONTINUE
C
      CALLCELL(Z,G)
C
      DO16L=1,10

```

```

C
      R(NR,L)=G(L)
16  CONTINUE
C
20  CONTINUE
C
C***** FORM TR MATRIX AND INVERT IT
C***** ****
C
IF (N69.EQ.69) GO TO 169
C
DO 30 I=1,10
DO 30 J=1,10
C
TR(I,J)=0.
C
DO 30 K=1,NR
C
TR(I,J)=TR(I,J)+R(K,I)*R(K,J)
C
30  CONTINUE
C
CALL MDINV
C
169  CONTINUE
C
C***** FORM U MATRIX
C***** ****
C
DO 35 I=1,10
DO 35 J=1,NR
C
U(I,J)=0.
C
DO 34 L=1,10
C
U(I,J)=U(I,J)+TR(I,L)*R(J,L)
C
34  CONTINUE
35  CONTINUE
C
C***** SUM OVER THE FITTING POINTS
C***** ****
C
DO 50 L=1,NR
C
C***** SET UP K VECTOR
C***** ****
C
ISUM=0
C
DO 41 K=1,3
C
IR(K)=IM(K,L)+ISVEC(K)+2
C
C***** CHECK ON LENGTH
C***** ****
C
ISUM=ISUM+IR(K)-1
C
41  CONTINUE
C

```

```

IDIF=ISUM-3*MESH
IF(IDIF.LE.0)GOTO043
C
C***** VECTOR OUTSIDE *****
C***** *****
C
DO42 I=1,3
C
    IR(I) = (MESH*2)-(IR(I)-1)
    IR(I)=IABS(IR(I))+1
C
42 CONTINUE
C
43 CONTINUE
C
C***** FOLD INTO PROPER ZINE: Y>X>Z *****
C***** *****
C
DO45 I=1,2
C
    I1=I+1
C
    DO45 J=I1,3
C
        IDIF=IR(I)-IR(J)
C
        IF(IDIF.GE.0)GOTO45
C
        ISAVE=IR(I)
        IR(I)=IR(J)
        IR(J)=ISAVE
C
45 CONTINUE
C
    IZ=IR(3)
    IX=IR(2)
    IY=IR(1)
C
C***** VECTOR NOW PROPER *****
C***** *****
C
INDEX=IP(IX,IY,IZ)
C
C***** SUM OVER THE BANDS *****
C***** *****
C
DO49 N=1,9
C
    ENG=EMESH(N, INDEX)
    PRS1=SME1(N, INDEX)
    PRS2=SME2(N, INDEX)
    PRD1=DME21(N, INDEX)
    PRD2=DME22(N, INDEX)
    PRD32=DME32(N, INDEX)
    PRD31=DME31(N, INDEX)
C
DO48 K=1,10
C
C***** SETUP FINAL SUM *****
C***** *****
C
CEP(N,K)=CEP(N,K)+U(K,L)*ENG

```

```

SEP1(N,K)=SEP1(N,K)+U(K,L)*PRS1
SEP2(N,K)=SEP2(N,K)+U(K,L)*PRS2
DE21(N,K)=DE 21(N,K)+U(K,L)*PRD1
DE22(N,K)=DE 22(N,K)+U(K,L)*PRD2
DE31(N,K)=DE 31(N,K)+U(K,L)*PRD31
DE32(N,K)=DE 32(N,K)+U(K,L)*PRD32

```

```

C
48 CONTINUE
49 CONTINUE
50 CONTINUE

```

```

C      N69 = 69

```

```

C ****
C      RETURN
C ****
C
END

```

```

C      SUBROUTINE SYM(S,N)

```

```

C ****
C      SUBROUTINE TO SYMMETRIZE THE HAMILTONIAN MATRIX
C      F.M. MUELLER MARCH 1966
C ****

```

```

C      DIMENSION S(36,36)

```

```

C      DO 10 I = 1,N
C      DO 10 J = I,N
10      S(J,I) = S(I,J)

```

```

C ****
C      RETURN
C ****

```

```

C
END

```

```

SUBROUTINE ZEP(CK,EZ)

```

```

C ****
C      J356 --FIND THE EIGEN VALUES FOR THIS POINT IN THE B.Z.
C      F.M. MUELLER MARCH 1966
C      MODIFIED BY S.G.DAS AUGUST 1969
C ****

```

```

C      DIMENSIONS MI1(9),DMI1(9),SMI(9),DMIT2(9),DMIT1(9),DMIT2(9),
C      DIMENSION DME21(9,300),DE21(9,10),SME1(9,300),SEP1(9,10),SME2(9,
8300),SEP2(9,10),DME31(9,300),DE31(9,10),DME32(9,300),DE32(9,
4300),DME22(9,300),DE22(9,10)
C      DIMENSION AL(20,9),DIFER(20)
1, CK(3),EZ(9)
C      DIMENSION PREM(16),BK(3),D(5),AD(12),S(36,36),EV(36),EG(9)

```

```

C
COMMON/GURU/DME21,DE21,SME1,SEP1,SME2,SEP2,DME31,DE31,DME32,DE32,
7DME22,DE22,SMI,SMI1,DMI1,DMI2,DMIT1,DMIT2
COMMON S,E,V,V0,V1,V2,V3,V4,ENGL,D,T,EL,R0,R1,EP,AD,AL,DIFER,BK
1,CNORM,MSIZE

```

```

C ****
C
14      FORMAT(2I6,F10.5)

```

00049712

00 05 00 20

```

DO9I=1,9
C
SMI=0.0
SMI1=0.0
DMI1=0.0
DMI2=0.0
DMIT2=0.0
DMIT1=0.0
C
9  CONTINUE
C
DO 20 I = 1,3
C
BK(I) = CK(I)
C
20  CONTINUE
C
DO30I=1,MSIZE
DO 30 J = 1,MSIZE
S(I,J)=0.
C
CALL HELD
CALL SYMS$5
CALL HSOC
CALL SYMS$9)
C
IF(MSIZE.LE.9) GO TO 35
CALL FILL
C
C*****PUT ON SMALL MAGNETIC FIELD TO SPLIT DEGENRACIES*****
C
C*****PUT ON SMALL MAGNETIC FIELD TO SPLIT DEGENRACIES***** 00 05 09 10
C
AMAG=0.0001
C
DO10I=1,9
C
J=I+9
L=I+18
M=J+18
C
S(I,I)=S(I,I)+AMAG
S(L,L)=S(L,L)+AMAG
S(J,J)=S(J,J)-AMAG
S(M,M)=S(M,M)-AMAG
C
10  CONTINUE
C
35  CONTINUE
C
ACC=.0001
C
CALL EIGENB(MSIZE,MSIZE,MSIZE,ACC)
C
MC=MSIZE/36
C
DO40I=1,9
C
I1=10-I
IT=(I1-1)*MC+3+I1
C
EZ(I) = EV(IT)+CNORM
EG(I) = EV(IT)+CNORM
C
40  DO11NLEVEL=1,9
C
ISUM=10-NLEVEL
C
00050070
00050992
00050993
00050994
00050995
00050996
00050970
00050980
00050990
00050991
00051400
00051910
00051920

```

```

C      DD1000 NGO=1,1
C
IS=ISUM+MC**3*( ISUM-1 )
IS=IS+NGO-1
C
D1=0.0
D2=0.0
D31=0.0
D32=0.0
S1=0.0
S2=0.0
C      DD28I=4,5
C
K=I+9
J=I+18
L=K+18
C
D1=D1+S(I,IS)**2+S(J,IS)**2
D2=D2+S(K,IS)**2+S(L,IS)**2
C      CONTINUE
C      DD28I=1,3
C
K=I+9
J=I+18
L=K+18
C
D31=D31+S(I,IS)**2+S(J,IS)**2
D32=D32+S(K,IS)**2+S(L,IS)**2
C      CONTINUE
C      DD271I=6,9
C
K=I+9
J=I+18
L=K+18
C
S1=S1+S(I,IS)+S(J,IS)
S2=S2+S(K,IS)+S(L,IS)
C      CONTINUE
C
S1=S1*S1
S2=S2*S2
C
SM1(NLEVEL)=S1
SM1(NLEVEL)=S2
DM1T1(NLEVEL)=D31
DM1T2(NLEVEL)=D32
DM1I1(NLEVEL)=D1
DM1I2(NLEVEL)=D2
C      CONTINUE
1000  CONTINUE
11    CONTINUE
C
C***** RETURN *****
C***** END *****

```

```
/*
//GO.SYSIN DD *
5
340
THE PARAMETERS FOR PLATINUM FOUND FROM THE WONDERFUL DHVA WORK OF LEES JOHN
0 .983
-1 .068900
0 .002700
0 .003300
0 .676799
-0 .496138
0 .011668
-0 .063561
0 .025417
0 .019276
-0 .000300
1 .949329
1 .792763
0 .370036
0 .268328
0 .024065
.6565
0 .974494
0 .001000
-0 .002300
0 .005400
-0 .001200
36
2200
1 .4
-0 .2
/*
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